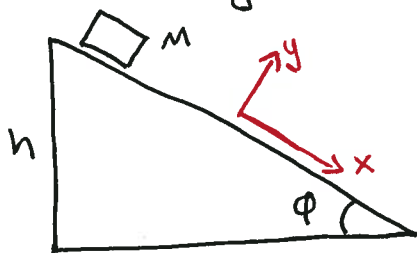


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force diagram example:



find an equation for the position of the block with respect to time

① force analysis, free body diagram

$$f_x = mg \cos \phi$$

② Newton's equation

$$f_x = m a$$

$$= m \ddot{x}$$

so

$$\ddot{x} = g \cos \phi$$

in physics, sometimes a dot is used to represent a time derivative:

$$v = \frac{dx}{dt} = \dot{x}$$

$$a = \frac{dv}{dt} = \frac{d}{dt} \dot{x} = \ddot{x}$$

③ Integration

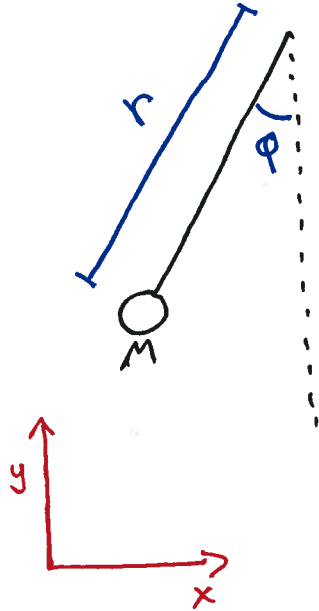
$$x = \iint \ddot{x} dt^2$$

$$= \frac{1}{2} g \cos \phi t^2$$

This is a very simple problem. for more complex systems we need a better approach.

Introducing more complex systems

1) Ball on string pendulum



- Two coordinate systems.
- Each has "two" dimensions
- Each has the same information

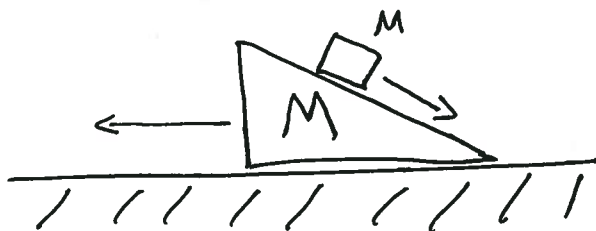
a) cartesian

b) polar

- Look closely: 1 dimensional system, ϕ !

Can you write down $x(t)$? $\phi(t)$?

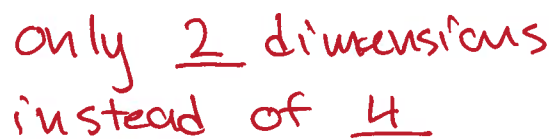
2) Wedge sliding under a block, no friction



- complicated!
- 4 position dimensions
- can you solve it?

Can't solve this easily

only the important quantities for energy and motion, etc.



B) Potential energy: mgh , $\frac{1}{2} k x^2$, ... (u)

$T = \text{kinetic}$; $U = \text{Potential}$

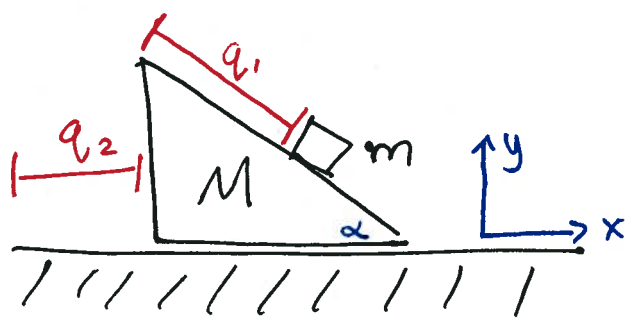
$$\mathcal{L} = T - U$$

And the equation

$$\frac{\partial \mathcal{L}}{\partial x} = \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{x}} \right)$$

give the equation of motion for a system. Let's look at an example

Example of Lagrangian



$$x_m = q_2 + q_1 \cos \alpha$$

$$y_m = -q_1 \sin \alpha$$

$$x_M = q_2$$

① what is the kinetic energy T ?

$$T_m = \frac{1}{2} m v^2 = \frac{1}{2} m (\dot{q}_1^2 + \dot{q}_2^2 + 2\dot{q}_1 \dot{q}_2 \cos \alpha)$$

$$(\text{This } v^2 \nearrow = \dot{x}_m^2 + \dot{y}_m^2)$$

$$T_M = \frac{1}{2} M \dot{q}_2^2$$

so $\textcircled{T} = T_m + T_M = \frac{1}{2} (M+m) \dot{q}_2^2 + \frac{1}{2} m (\dot{q}_1^2 + 2\dot{q}_1 \dot{q}_2 \cos \alpha)$

② what is the potential energy?

$$\textcircled{U} = mgy_m$$

$$= -mgy_m = -mgq_1 \sin \alpha$$

③ write down the Lagrangian $\mathcal{L}$

$$\mathcal{L} = T - U$$

$$= \frac{1}{2} (M+m) \dot{q}_2^2 + \frac{1}{2} m (\dot{q}_1^2 + 2\dot{q}_1 \dot{q}_2 \cos \alpha) + mgq_1 \sin \alpha$$

recall: $\frac{\partial \mathcal{L}}{\partial x} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{x}}$

In this example, we have 2 coordinates

$$x \rightarrow q_1, q_2$$

So we have 2 equations

$$\left\{ \begin{array}{l} \frac{\partial \mathcal{L}}{\partial q_1} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}_1} \\ \frac{\partial \mathcal{L}}{\partial q_2} = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}_2} \end{array} \right\}$$

don't get intimidated by these, they are usually very simple equations

Let's solve the q_1 equation first. Remember, these are partial derivatives, so $\frac{\partial q_2}{\partial \dot{q}_1} = \frac{\partial \dot{q}_2}{\partial \dot{q}_1} = 0!$

$$(\mathcal{L} = \frac{1}{2}(M+m)\dot{q}_2^2 + \frac{1}{2}m(\dot{q}_1^2 + 2\dot{q}_1\dot{q}_2\cos\alpha) - mgq_1\sin\alpha)$$

Plug in:

$$mg\sin\alpha = \frac{d}{dt}(m[\dot{q}_1 + \dot{q}_2\cos\alpha])$$

$$\boxed{mg\sin\alpha = m\ddot{q}_1 + m\ddot{q}_2\cos\alpha}$$

The other eqn for q_2 gives

$$\boxed{\ddot{q}_2(m+M) = m\ddot{q}_1\cos\alpha}$$

Algebra: $\ddot{q}_1 = \frac{g\sin\alpha}{1 - \frac{m\cos^2\alpha}{M+m}}$

This is all we need to know about the system!

- 1) write down the lagrangian for a ball in free fall

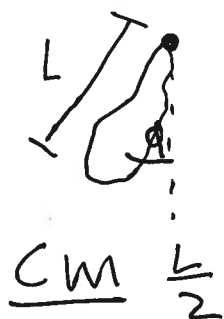
$$\mathcal{L} = \frac{1}{2} m \dot{x}^2 - mgx$$

- 2) write down the lagrangian ^{for system} we looked at earlier:



$$\mathcal{L} = \frac{1}{2} m \dot{q}^2 - mgsin\alpha$$

- 3) Tricky: write Lagrangian for a pendulum with moment of inertia I :



$$\mathcal{L} = \frac{1}{2} I \dot{\phi}^2 + mg \frac{L}{2} \cos \phi$$

Symmetry & Lagrangian

Symmetry can tell you a lot

"I have a rotationally symmetric object"



It's a ball

Symmetry in mechanics:

- Rotational $\longleftrightarrow$ conservation of angular momentum
- Translation $\longleftrightarrow$ conserve momentum (linear)
- Time $\longleftrightarrow$ conserve energy

To say that the Lagrangian $\mathcal{L}$ is rotationally symmetric is to say

$$\frac{\partial \mathcal{L}}{\partial \phi} = 0$$

This implies certain constraints on what $\mathcal{L}$ can be.

By picking symmetries, you determine the form of the Lagrangian.

Then the Lagrangian determines the behavior of the system.

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The standard model Lagrangian

Symmetry:

4d Poincaré &

$$SU(3)_C \times SU(2)_L \times U(1)_Y$$

format:

$$\mathcal{L} = \mathcal{L}_{\text{kinetic}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{Yukawa}}^{(\text{flavor, strong})}$$

This is the most general form of Lagrangian given above symmetries

There is a lot of detail hidden in ~~this~~ this Lagrangian (which I don't understand)

However, this is a Lagrangian like the ones we found before.

This Lagrangian specifies particle masses, interactions, etc. It describes

everything we know about
particle physics.